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2 **Electrochemical Probe**

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14 **Abstract**

15 Electrochemical sensors can be used to create portable, attractive, high-sensitivity
16 devices to detect heavy ions in water at a reasonable cost. A novel method of
17 measuring local redox on an electrode using the differential reflectivity measured on
18 about 6 micron diameter spot using a He/Ne laser beam is described. Field focusing
19 and electrode modification enhances the sensitivity, making the method ideal for
20 miniaturization and multiplexing multiple analytes on a monolith electrode. The method
21 is demonstrated by detecting Pb, Hg, and As ions at ppb and ppt levels.

22

23 **Keywords:** Heavy metals; electrochemical analysis; electrode size; dielectrophoretic
24 effect; water pollution; electro-optics

25

26 **1. Introduction**

27 With ever-decreasing sources of water for the growing world population, the search for
28 new sources and the recycling of water have made the issue of heavy metal ion
29 contamination a ubiquitous problem and intensified the need for new detection systems,
30 in terms of sensitivity, specificity, and portability. Toxicity of heavy metal ions is well
31 documented [1-3] with several toxic metals, such as As, Cd, Pb, and Hg, having no
32 known biological function [4]. Heavy metal ions are toxic to reproductive organs of
33 females [5] and males [6], and some are also carcinogens [7]. Concentrations of Pb,
34 Hg, and As in ppb levels a normal human body pose serious health hazards, such as
35 plumbism, hypertension, emphysema, and even cancer with chronic, irreversible
36 damage to organs such as kidneys, skeletal muscles, and soft tissues [8]. Highly
37 sophisticated methods, such as inductively coupled plasma-mass spectroscopy (ICP-
38 MS) have been developed to measure Cd, Hg, and Pb ions at a sensitivity of 4, 6, and
39 13 ppb, respectively [9]. Optical methods, such as atomic absorption spectroscopy
40 (AAS), X-ray fluorescence spectroscopy (XRF), and inductively coupled plasma-optical
41 emission spectroscopy (ICP-OES), have been developed to measure heavy metal ions
42 at ppb levels [10, 11].

43 Arguably, electrochemical sensors are the most pervasive [11] with innovations in
44 paper-based disposable devices [12] and enhanced sensitivity using nanomaterials [13-
45 15] and biopolymers [16]. Apart from the simplicity and expense, electrochemistry is an
46 attractive detection principle because the (redox current) signal is linearly proportional to
47 the analyte concentration. However, the signal (and hence the sensitivity) diminishes as
48 the size of the sensor is reduced. The modulation of the refractive index due to the
49 change in ionic interface (i.e., electrical double layer (EDL)) during electrochemical
50 reaction has also been leveraged by static [17] and dynamic [18] interferometry and

51 surface plasmon resonance [19,20]. Here, a simple optical technology based on
52 differential reflectivity is described where the signal increases as the sensing area
53 decreases to naturally decrease the device sensing size to micron levels. The principle
54 of the method and its application for detecting Pb, Hg, and As ions at ppb levels is
55 described.

56 2. Results and Discussion

57 The electrochemical reaction (of the redox moiety) was quantitatively probed by
58 measuring the modulation in reflectivity of the sensing electrode due to the change in
59 the interfacial ionic environment around the redox potential (Fig. 1(a)). The
60 electrochemical system consisted of a standard three-electrode arrangement used in
61 cyclic voltammetry (CV). The sensor's working electrode (WE) was made of bare Au in
62 contact with a redox ion containing an aqueous solution. The reference electrode (RE)
63 was Ag/AgCl, and Pt wire served as a counter electrode (CE). A potential, E, was
64 applied between the RE and WE; and the current, I, was measured between WE and
65 CE using a potentiostat (Autolab PGSTAT 128N). The WE was patterned with SU8
66 photoresist to expose a circular spot of diameter, d, ranging from 10 to 300 μm .
67 Reflectivity of a He-Ne laser with a beam diameter of about 6 μm was measured from
68 the WE during the CV. The reflectivity was measured as a differential signal by applying
69 additional AC potential of frequency, $\omega = 0.5$ kHz, on the WE. The reflected light
70 intensity, R_A , modulated at ω due to ion oscillation caused by AC potential similar to that
71 observed by differential interferometry [18, 21]. From the intensity of the constant
72 reflected light, R_O , measured as a DC signal, differential reflectivity was obtained as, R
73 $= R_A/R_O$. A typical signal for R as a function of multiple CV cycles from -0.1 to 0.6 V at a
74 ramp rate of 0.5 V/s for 50 mM of $\text{K}_4[\text{Fe}(\text{CN})_6]$ showed a highly periodic signal with
75 excellent cycle-to-cycle reproducibility. About 75 scans in 3 minutes were obtained for
76 analysis. For better visualization, magnified view of only a part of the data is shown to
77 underscore the periodicity (Fig. 1(b)). Overlaying all of the cycles and filtering high
78 frequency noise, the average signal, $\langle R \rangle$, versus E with robust statistics, was obtained
79 (Fig. 1(c)). The error-halo around $\langle R \rangle$ was the scatter due to cumulative cycle-to-cycle
80 variations. Typically, the noise level at 500 Hz is less than $10^{-8}R_O$. For no ions (i.e.,
81 blank), $\langle R \rangle$ is nominally a flat line implying that the signal is absolute.

82 The physical principle of the measurement is considered to relate R to (bulk) ion
83 concentration, [c]. For normal beam, the reflectivity, r by Fresnel's law is, $r = |(N_1 - N_2) /$
84 $(N_1 + N_2)|$, where N_1 and $N_2 (=n - ik)$ are complex refractive index of the solution at the
85 interface and metal, respectively. As the ion concentration in the EDL changes due to
86 applied potential, $N_1 (=n_1 - ik_1)$ will modulate leading to change in r. Due to small AC
87 potential at ω , the ion concentration at the interface will also oscillate at ω with an
88 amplitude of δc_0 . (In principle, ion modulation, δc will decrease as distance from
89 electrode increases. For estimation, with no loss in formulation of the basic concept of
90 the measurement, spatially averaged modulation, δc_0 is considered). As differential
91 refractive index, dn/dc is constant, for small AC potential, $n_1 = N + (dn/dc)\delta c_0 \cos(\omega t)$,
92 where, N is the refractive index at given E (i.e., equilibrium). The higher order oscillation
93 of R_A at 2ω to 6ω measured by the lock-in amplifier is $>10^3$ -fold smaller than R_A at ω .
94 (The higher order effects may occur because of non-linear change in the ion
95 concentration gradient as a function of distance from the electrode/solution interface.)

96 Assuming, the solution to be non-absorbing (i.e., $k_1 = 0$) and for Au, $n^2 \ll n_1^2$; and
97 because the nonlinearity can be neglected (at 0.1% accuracy), the Fresnel's law for
98 dynamic n_1 can be linearized to yield, $|r|$ (at ω), $R = Q[(dn/dc)\delta c_0]\cos\omega t$, where $Q =$
99 $[2n(N^2-k^2)/(N^2+k^2)]$.

100 Next relationship between E and δc_0 is considered to relate the observed R_{\max} to $[c]$. It
101 is well known that, due to slow diffusion kinetics of ions during electrochemical reaction,
102 there is a significant deviation from equilibrium of the ion concentration profile at the
103 interface over a few microns [22]. As a result, the charge in the EDL significantly
104 deviates from neutrality leading to insufficient screening causing a deeper penetration of
105 the electric field emanating from the electrode. This maximization of oscillation of ion
106 amplitude close to redox potentials due to the AC field is consistent with independent
107 (and explicit) observation by differential interferometry [18,21]. Therefore, the reflectivity
108 due to oxidation and reduction peaks at ${}^oR_{\max}$ and ${}^rR_{\max}$, respectively (Fig. 1(c)). At E far
109 from the redox potential, δc_0 is small leading to the small baseline (Fig. 1(c)).
110 Furthermore, as ion diffusion is significantly slower than electron transport in the
111 electrode, EDL discharge is linearly proportional to redox current which in turn is linearly
112 proportional to bulk ion concentration, $[c]$. Thus, δc_0 is linearly proportional $[c]$. Thus,
113 after baseline subtraction, at the redox, $R_{\max} = PQ[c]$, where P is a calibration constant
114 that will depend on the dynamic properties of the target ions, such as diffusion constant
115 and electrophoretic mobility that determines δc_0 .

116 By scanning the laser beam and measuring R_{\max} , local binding on a monolith electrode
117 can be mapped. If the electrode were to be fabricated as a microarray of a local surface
118 modification, multianalyte detection can be obtained. From known P , the target $[c]$ can
119 be quantified. Importantly, as the polarizability of the ion will change significantly due to
120 its charge state of the ion, a large modulation in n_1 is expected, leading to high
121 responsivity; while k_1 is nominally constant (as change in **absorption** at 633 nm is not
122 large).

123 A typical $\langle R \rangle$ versus E for 50 mM showed a large variation in ${}^o/rR_{\max}$ as a function of the
124 patterned spot size, d . As shown by differential interferometry [21], the signal increased
125 as d decreased due to field focusing (i.e., dielectrophoresis) (Fig. 2(a)). The signal was
126 further enhanced as the thickness of the photoresist increased from 500 to 1200 nm
127 (compare $d = 50 \mu\text{m}$ for 500 and 1200 nm thickness in Fig. 2(a)). Conventional CV (not
128 shown) indicated, as expected, the opposite effect where the peak oxidation current
129 decreased from 50 μA to 20 μA as the size of the spot decreased from 300 μm to 50 μm
130 due to a smaller electrode surface area. A systematic change in d showed the classic
131 field focusing effect where the current density scales as $1/d$ (Fig. 2(b), inset) [23].
132 Because, the laser beam diameter was fixed, R_{\max} corresponded to the current density
133 rather than the total current [21]. The linear fit for $1/d$ versus R_{\max} was excellent with a
134 small intercept. The intercept of 0.46 corresponds to (limiting) signal for a large planar
135 electrode. Thus, for a spot size of 50 μm and a thickness of 1200 nm, the enhancement
136 due to field focusing relative to a large electrode was ~ 13 -fold. For a larger thickness
137 (i.e., 1200 nm in Fig. 2(b)), the signal began to decrease at a smaller d because of the
138 incomplete etching during development of SU8 and the limitations of the lithography
139 process (i.e., dishing effect). For a fixed $[c]$ of 20 ppb, a similar size effect as $[\text{Fe}(\text{CN})_6]^{4-}$
140 $^{3-}$ is observed with a robust signal for Pb, Hg, and As for $d = 50 \mu\text{m}$. (Fig 2(c), data only)

141 from the oxidation peak is shown). The signal is nominally enhanced by ~5-fold as the
142 spot size is reduced from 300 to 50 μm at fixed thickness of 1200 nm.

143 The dielectrophoretic effect was leveraged to measure heavy metal ions (Fig. 3). The
144 three ions tested were As (in As_2O_3), Hg (in HgCl_2), and Pb (in PbCl_2). The calibration
145 curve for all the ions studied is highly linear with fitness of >0.95 . The sensitivity to
146 detect Pb ion is significantly enhanced by 3-folds by modifying the Au surface with
147 adsorbed poly(styrene sulfonate) (PSS) (Fig. 3(a)). A robust peak is obtained at 50 ppt
148 (Fig. 3(a), inset). The peak for a (baseline corrected) R_{max} of 0.17×10^{-4} is difficult to
149 discern, leading to estimated limit of detection (LOD) of ~15 ppt. At ppb levels, the
150 sensitivity of As ion in ppb levels is 1.5 levels lower than Pb on bare Au (Fig. 3(b)). By
151 modifying the Au electrode with PSS and imbedding $[\text{Ru}(\text{NH}_3)_6]^{3+}$ ions using 32 cycle of
152 CV ranging from -0.5 to 1.2 V, similar to a method described earlier [24], a robust peak
153 is obtained at 1 ppt (Fig. 3(b), inset). Peak below 1 ppt is difficult to discern, leading to
154 LOD of ~1 ppt. Sensitivity to detect Hg ion on bare Au was about 1.5 fold higher than Pb
155 ion at ppb levels (Fig. 3(c)). By imbedding $[\text{Fe}(\text{CN})_6]^{4-}$ ions on modified Au with PSS and
156 poly(allylamine hydrochloride) PAH layers using the method described earlier [24], a
157 robust peak is obtained at 1 ppt (Fig. 3(c), inset). For Hg, LOD was ~0.5 ppt
158 corresponding to R_{max} of $\sim 0.068 \times 10^{-4}$ (close to 1 ppt). All the ion detection was
159 performed at $d = 50 \mu\text{m}$ and resist thickness of 1200 nm. As indicated for Fig. 1(c), the
160 signal with no ions (i.e., blank) for each modified electrode in the insets of Fig. 3 was
161 zero indicating that R_{max} is an absolute measurement. The mechanism of the imbedded
162 redox ion as described earlier [24], is to function as electron mediator to enhance the
163 redox signal and thereby improve sensitivity. The error bar is calculated from the
164 magnitude of the error halo at R_{max} . As the magnitude of the signal reduces, the width of
165 the error halo also nominally reduces, however the relative error at low $[c]$ is high as
166 shown in the semilog plot (Fig. 3(d)). To demonstrate the quality of the signal averaged
167 over 75 cycles, the inset of Fig. 3(a) to 3(c) shows the raw R with an error halo at the
168 lowest $[c]$ for each ion with no baseline correction. The potentials for R_{max} are consistent
169 with the formal potentials of redox for As, Hg, and Pb ions which are at 320, 380, and
170 280 mV, respectively. Permissible safe levels of As, Hg, and Pb in drinking water set by
171 the United States Environmental Protection Agency are 10, 10, and 15 ppb,
172 respectively. Thus, the low LOD and strong signal at ppb levels in water demonstrates
173 the feasibility of this approach to measuring the toxicity of heavy metal ions in water.
174 Importantly, as multiple redox reactions can be measured on a monolith electrode, by
175 locally modifying the electrodes, it will potentially be possible detect complex mixtures of
176 heavy ions. The reasonableness of the signal (in Fig. 3) is inferred by estimating the ion
177 modulation for measured R . To first order, the dn/dc will primarily be due to the anion
178 owing to its (much) higher polarizability than the cation. Thus, for the chloride salts used
179 in the study, $dn/dc \sim 10^{-1}/\text{M}$ for Hg and Pb salts (based on NaCl). At equilibrium, $N \approx 1.4$
180 (the ions will raise the interfacial refractive index of water by ~10%). For Au ($n = 0.18$
181 and $k = 3.07$ at 633 nm) electrode, far from redox, $R \sim 0.5 \times 10^{-5}$ leads to $\delta c_0 \sim 2.5 \text{ mM}$.
182 For 500 ppb solution, $R \sim 10^{-4}$ at redox corresponds to δc_0 of ~50 mM. Although the rise
183 in concentration due to AC potential seems high, for $[c]$ of 500 ppb (i.e., ~2 μM), the
184 enhancement of $[\text{Cl}]^-$ in EDL(at equilibrium) at 200 to 300 mV electrode potential would
185 be in 10^0 to 10^2 mM range from classical Poisson-Boltzmann distribution theory. Thus,
186 the rough estimation of δc_0 is reasonable within an order of magnitude. Furthermore,

187 consistent with the observation, for polyelectrolyte modified electrode n will increase
188 and k will be lowered causing an increase of $[2n(N^2-k^2)/(N^2+k^2)]$ to make the electrode
189 more sensitive.

190 191 **3. Conclusions**

192 In summary, the differential reflectivity method to measure current density is described
193 where the signal enhances significantly due to field focusing effects and electrode
194 modification. The signal is absolute. Limit of detection of 1 ppt is demonstrated for Hg
195 and As ions, and 50 ppt for Pb ions. The calibration curve for all the three ions was
196 highly linear over $[c]$ ranging over five orders of magnitude. Furthermore, as the method
197 measures local redox, potentially, multiple analytes on a monolith electrode can be
198 detected by, for example, locally modifying the electrode surface with different
199 polyelectrolytes and redox ions to develop a quantitative "artificial nose" to profile heavy
200 metal ions.

201 202 203 **References**

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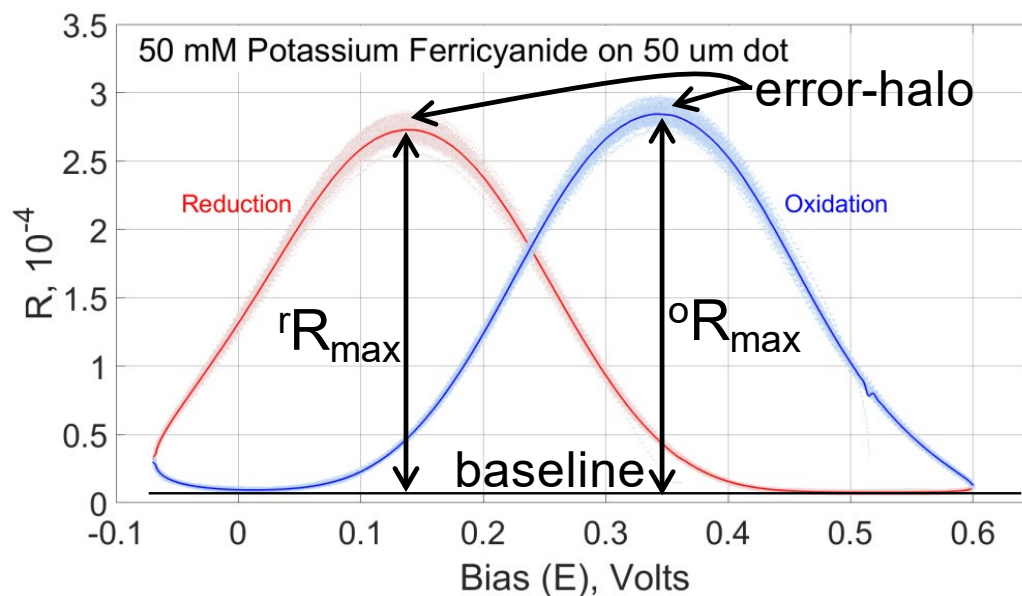
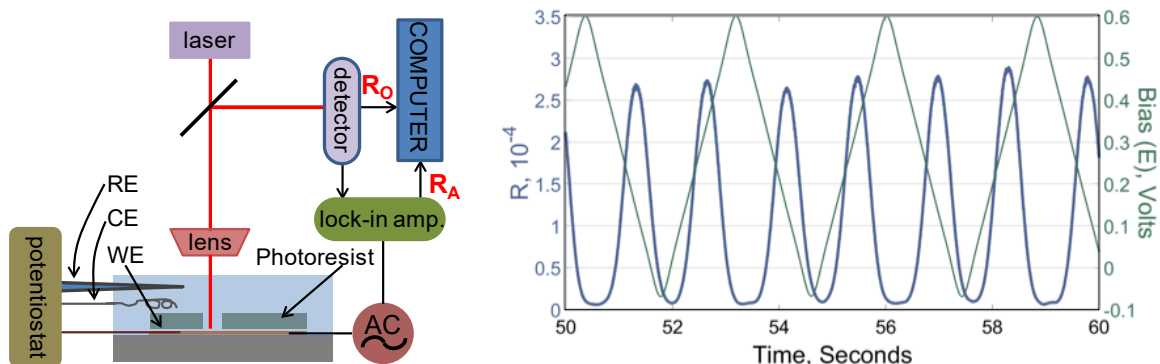


Figure 1. (a) Schematic displaying the optical set-up of the differential reflectometer. (b) Magnified scan of the raw data of reflectivity as a function of time showing highly repeatable redox peaks of $[\text{Fe}(\text{CN})_6]^{4-/3-}$. (c) The $\langle R \rangle$ as a function of E by superposing about 75 scans from the raw data.

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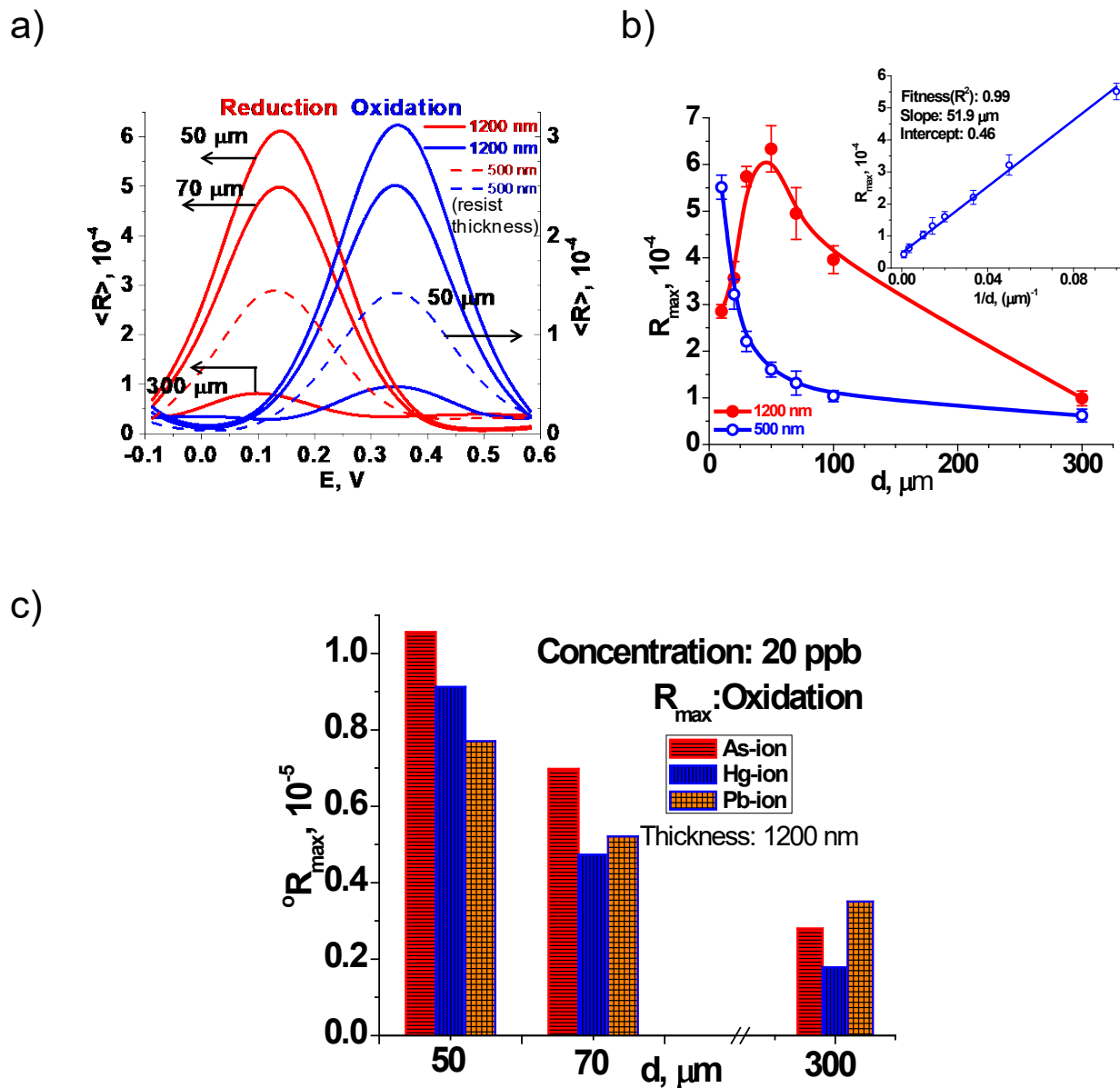
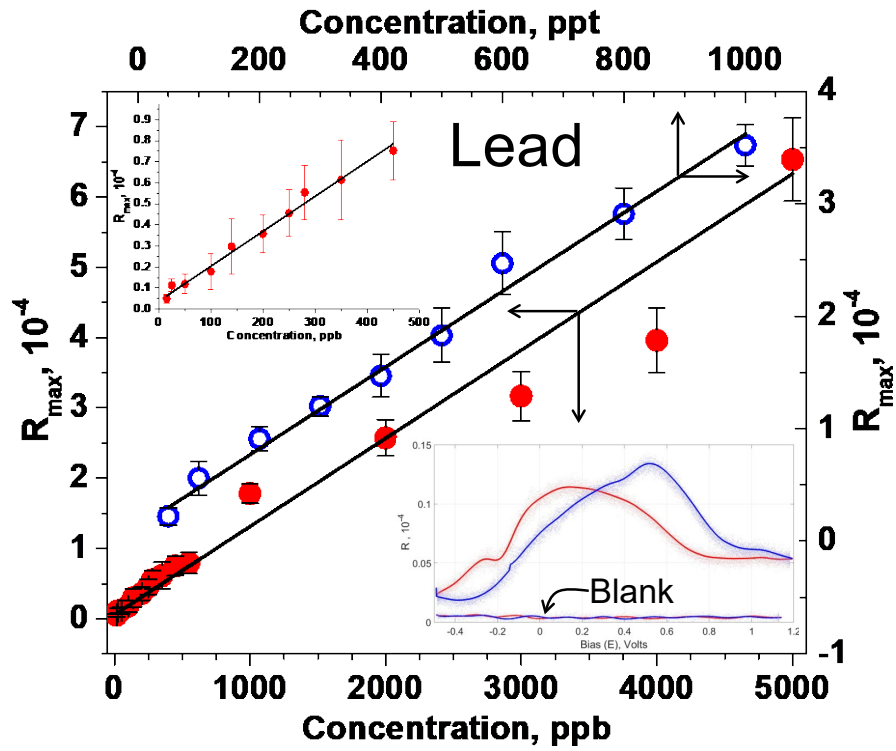
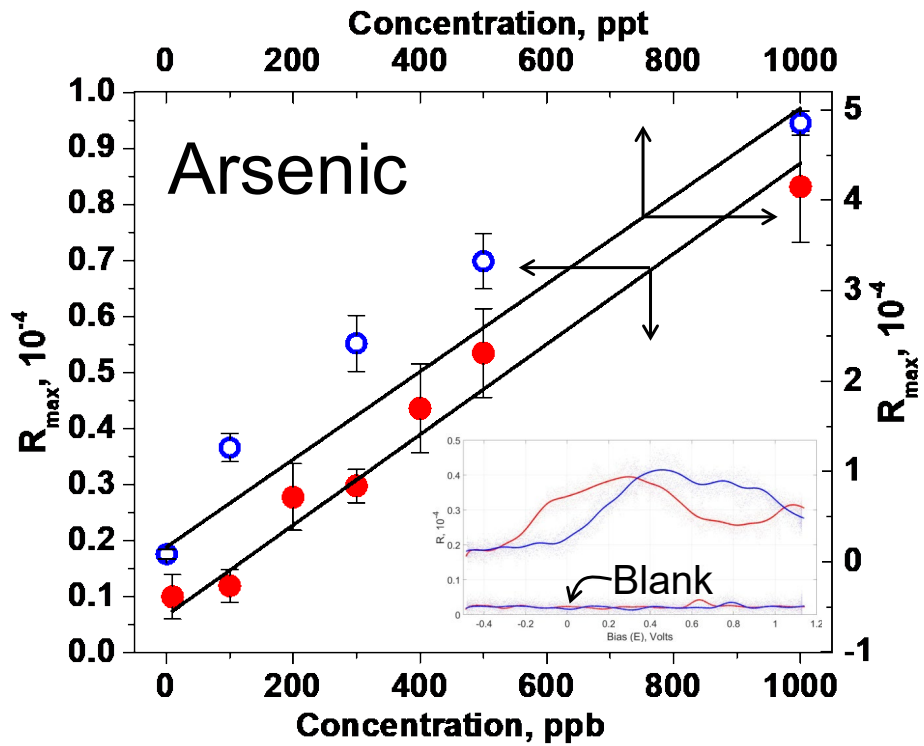


Figure 2. (a) Redox signal for 50 mM $K_4[Fe(CN)_6]$ as a function of resist thickness (nm) and the electrode diameter (μm). (b) R_{max} as a function of d for resist thickness of 1200 and 500 nm. Inset: Linear correlation of R_{max} versus $1/d$ for resist thickness of 500 nm. (c) The variation in R_{max} for heavy metal ions with d .

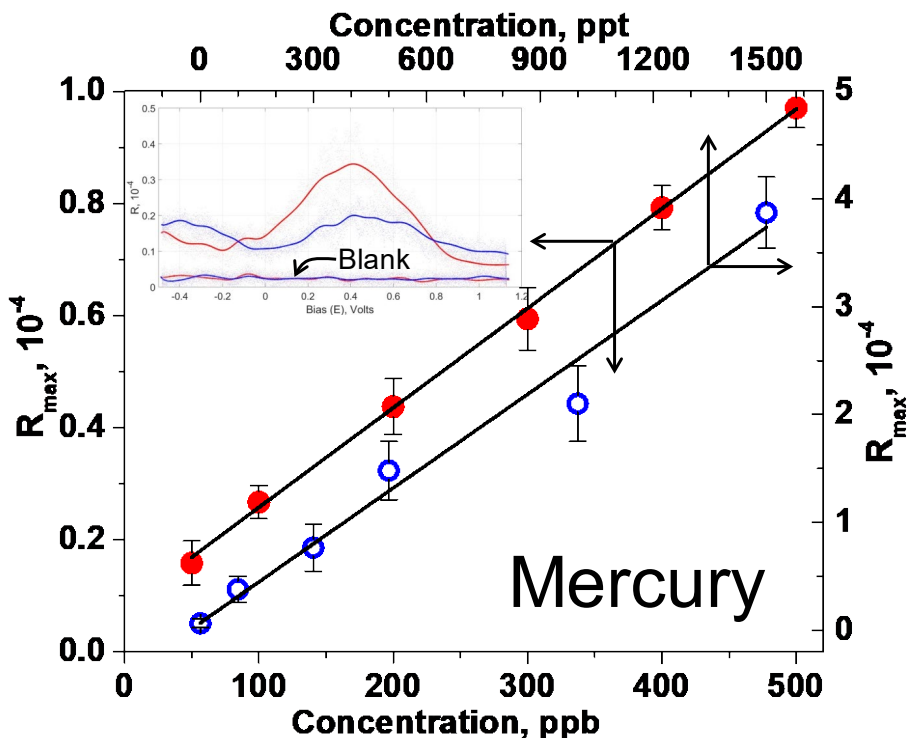
323 a)



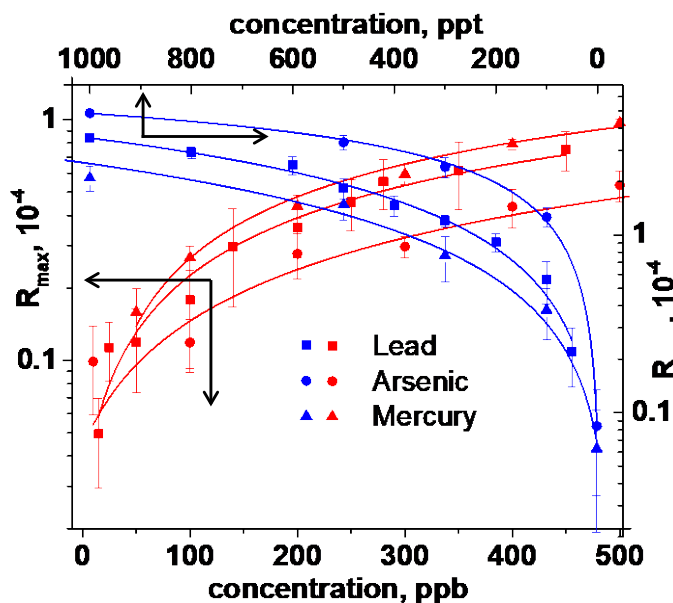
339 b)



355 c)



367 d)



379 **Figure 3.** The calibration curve of heavy metal ions with linear relation between R_{max}
 380 and $[c]$ ($d= 50$ and resist thickness = 1200 nm). (a) Pb on bare Au (ppb levels) and PSS
 381 modified Au (ppt levels). The responsivity (i.e., slopes) are 0.00125 and 0.00350 for the
 382 ppb and ppt levels, respectively. The inset shows raw data for 50 ppt Pb. (b) As on bare
 383 Au (ppb levels) and PSS modified Au with imbedded $[Ru(NH_3)_6]^{3+}$ ions (ppt levels). The
 384 responsivity are 0.00081 and 0.00486 for the ppb and ppt levels, respectively. The inset
 385 shows raw data for 1 ppt As. (c) Hg on bare Au microelectrode (ppb levels) and PSS
 386 and PAH modified with $[Fe(CN)_6]^{4-}$ ions (ppt levels). The responsivity are 0.00178 and
 387 0.00245 for the ppb and ppt levels, respectively. The inset shows raw data for 1 ppt Hg.
 388 The blank in the insets is R to the corresponding modified electrode in distill water. (d)
 389 Semilog plot of calibration curves to show relative error.